

QUARK-GLUON PLASMA  
AND  
RELATIVISTIC ION COLLISIONS

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November 1984

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AND  
RELATIVISTIC ION COLLISIONS.

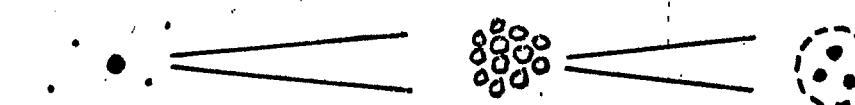
L. VAN HOVE

A labwide lecture and two seminars  
at Brookhaven National Laboratory  
(November 1984)

Atom

Nucleus

Nucleon



electron  
(lepton family)

nucleons  
(hadron family)

quarks  
(quark family)

Strong interaction: acts on quarks,  
is mediated by gluons.

Electro-Weak interactions: act on quarks and leptons,  
are mediated by photons and weak bosons.  
 $(W^\pm, Z^0)$

Leptons, photons, weak bosons, hadrons can exist as isolated particles.  
Quarks and gluons don't (confinement property).

## Confinement - Hadronization

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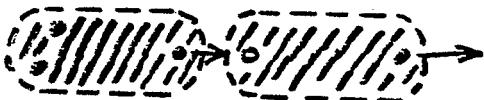
$$1 \text{ fm} \\ = 10^{-13} \text{ cm}$$



||||||| Colour electric field



QCD flux tube (also called string).



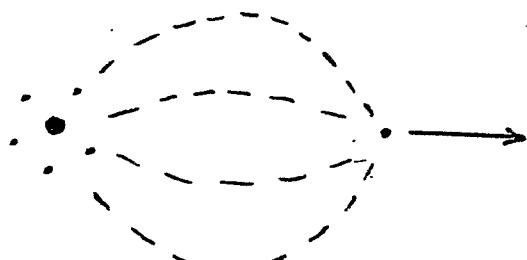
Hadronization by  $q\bar{q}$  pair creation.



String tension  $\sim 1 \text{ GeV/fm}$

For comparison, "liberation" of electron:

$$1 \text{ \AA} \\ = 10^{-8} \text{ cm}$$



Electric field

## Theoretical prediction of a new state of matter,

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## the "QUARK-GLUON PLASMA" (also called "quark matter").

- Constitution of hadronic matter according to

QUANTUM CHROMODYNAMICS (QCD),

a non-abelian gauge theory based on  $SU_3$  (colour qu. nbs)

[Spinor fields  $\rightarrow$  QUARKS,  $SU_3$  triplets, 5 (6!) "flavours"  
gauge fields  $\rightarrow$  GLUONS,  $SU_3$  octets, no further int. qu.n.

- CONFINED property: hadronic systems in vacuum form  $SU_3$  singlets ("colourless")

Examples: baryon =  $(qqq)_{\text{singlet}}$ , meson =  $(q\bar{q})_{\text{singlet}}$ ,



where  $q, \bar{q}$  may be "dressed" by a "sea" of  $q\bar{q}$  pairs and gluons.

- nucleus is also  $SU_3$  singlet, containing 3A quarks and perhaps more ( $q\bar{q}$  pairs).

-  $q, \bar{q}$  distribution in nuclei is not the one in nucleons folded with Fermi distribution (EMC effect).

- At low density, hadronic matter forms HADRON GAS.

- Single hadrons occupy finite volume,  $\approx 1 \text{ fm}^3$ .

- What about very dense hadronic matter as must occur if when nuclear matter is highly compressed (dense stellar objects) ii) at high temperature, high density of hadrons being created by thermal agitation (early Universe).

Hadrons must then coalesce into a dense, continuous fluid of  $q$ 's,  $\bar{q}$ 's and gluons, the QUARK-GLUON PLASMA.

i) Compress nuclear matter by factor 20:

Quark density  $\sim 20 \times 3 \times 0.17 \text{ fm}^{-3} \sim 10 \text{ fm}^{-3}$

ii) Heat matter to  $T \sim 500 \text{ MeV} \sim 6 \times 10^{12} \text{ K}$ :

An ideal pion gas would have  $\sim 6$  pions/ $\text{fm}^3$ ,  
i.e., a  $q + \bar{q}$  density  $\sim 12 \text{ fm}^{-3}$

Under such conditions, existence of Quark-Gluon Plasma  
seems inescapable.

- Can this new state of matter be produced on earth?  
Perhaps in ultrarelativistic nuclear collisions

- 3 classes of problems, all affected by  
large uncertainties:

i) Expected properties of QGP and of QGP  $\rightarrow$  HG transition,  
(HG = Hadron Gas).

ii) Compression and heating in nuclear collisions.

iii) Signals for QGP formation in " " .

- Much theoretical work, great quantitative uncertainties.

- Severe lack of data hampers work; experimental programmes to start soon at

Brookhaven (15 GeV/nucleon up to  $^{32}\text{S}$ ) and

CERN (225 GeV/nucleon up to  $^{16}\text{O}$ ).

- Impossibility to measure "colour conductivity".

RHIC: colliding beams up to gold ( $^{197}\text{Au}$ )

up to 100 GeV/nucleon.

(RHIC: Relativistic Heavy Ion Collider)

# Quark-gluon plasma (QGP) and hadronic phase transition

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One expects 2 phase transitions, but they may coincide:

Deconfinement transition:  $QGP \leftrightarrow HG$  (hadron gas)

Chiral transition: restoration of chiral symmetry  
(from  $\langle \bar{q} q \rangle \neq 0$  to  $\langle \bar{q} q \rangle = 0$ )

Theoretical work follows 2 approaches:

Continuum approach:

describe HG from hadron phenomenology

" QGP by ideal gas with perturbative and "plasmon corrections" based on QCD.  
(complexity of collective effects!)

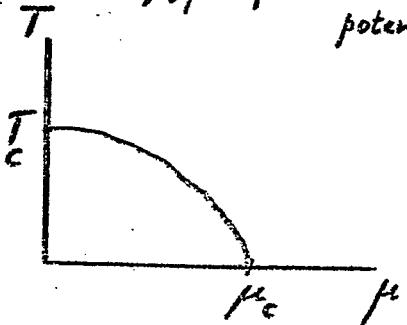
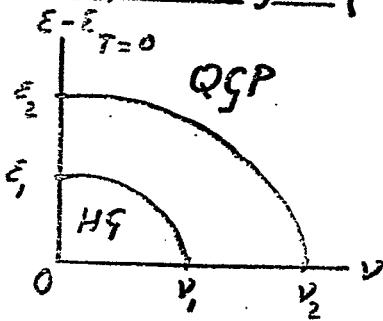
Lattice approach: thermodynamics of QCD in lattice approximation by Monte Carlo integration.

- Pure  $SU_3$  gauge field (self-interacting gluons):  
first order phase transition, large latent heat -

$(E_{QGP} \gg E_{HG}$  at transition,  $E$  = energy density )

- Full QCD with quarks } nature of transition unclear,  
but large variation of  $E$  for small variation of  $T, \mu$ .

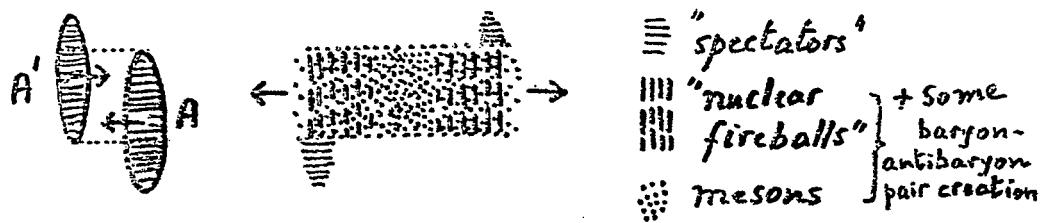
Possible phase diagram ( $\nu$  = net quark density,  $\mu$  = quark chemical potential)



perhaps  $\begin{cases} E \sim 2E_1 \sim \text{a few GeV fm}^{-3} \\ \nu_2 \sim 2\nu_1 \sim \text{a few fm}^{-3} \end{cases}$  | perhaps  $\begin{cases} T_c \sim 150-300 \text{ MeV} \\ \mu_c \sim 400-800 \text{ MeV} \end{cases}$   
(1 MeV =  $1.16 \times 10^{10} \text{ K}$ )

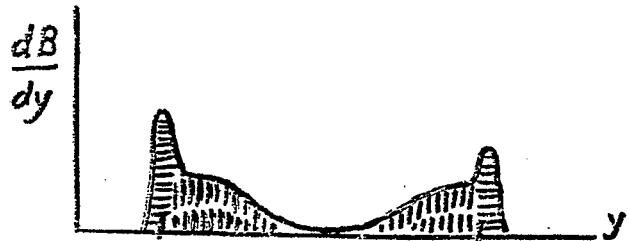
# Ultrarelativistic nuclear collisions.

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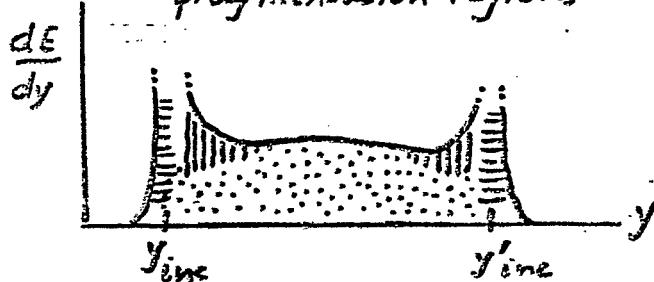
Longitudinal rapidity:  $y = \operatorname{arctanh} v_{\text{long.}}$

Net baryon number:  $B$



$y_{\text{min}} \longleftrightarrow y'_{\text{min}} = y_{\text{min}} + \Delta$   
central region

$\xrightarrow{\sim 2}$  fragmentation regions  $\xleftarrow{\sim 2}$



( $\Delta = 6$  for  
 $E_{\text{lab}} = 200 \text{ GeV}/N$ )

$dE$ : comoving energy of particles in  $(y, y + \Delta y)$

Fragmentation regions: compression gives  $v \sim 2y \sim 1 \text{ fm}^{-1}$ ,  
 $\epsilon \sim 1 \text{ GeV fm}^{-3}$  (maximum at  $E_{\text{lab}} \sim 15-20 \text{ GeV}/N$  ?)

Central region: perhaps  $\epsilon \gtrsim 2 \text{ GeV fm}^{-3}$  at  $E_{\text{lab}} \gtrsim 200 \text{ GeV}/N$   
(keeps growing with energy)  
But: FLUCTUATIONS can help!

## Signals for plasma formation.

### "Traditional" signals:

- direct  $e^+e^-$  and  $\mu^+\mu^-$  pair production
- direct photon production:

  	<p>to be calculated in plasma as function of <math>T</math> and <math>v</math>, etc. space-time distribution of <math>T, v</math> needed, production in absence of plasma poorly known.</p> <p>- production of heavy quark flavours (strange particles, charmed particles too heavy?): not in equilibrium in plasma, mainly produced in hot plasma, probably more abundant than without plasma (production by gluons mainly).</p>
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### Signals of phase transition with large latent heat:

- $\langle p_t \rangle$  - multiplicity correlation (observed in central region for  $p_T$  at  $E_{cm} \sim 50-60 \text{ GeV}$ ,  $\bar{p}p$  at  $E_{cm} \sim 540 \text{ GeV}$ ):  
Characteristic of thermal behaviour:

If plasma formation is common

$\langle p_t \rangle$  grows with temperature  $T$  and pressure  $p$ ,  
multiplicity grows with entropy.

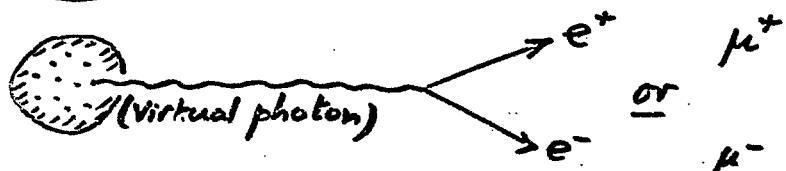
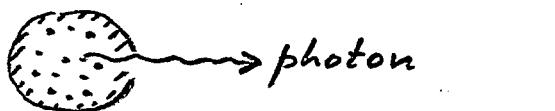
But phase transition with large latent heat gives large increase of entropy without increase of  $T, p$ .  
This might show as an anomaly in  $\langle p_t \rangle$  - mult. curve.

- If plasma formation is exceptional:
- multi-hadron fluctuations caused by instabilities in release of latent heat by plasma (plasma deflagration)  $QGP \xrightarrow{\text{Shock front}} HG \rightarrow$   
Cosmic ray evidence?

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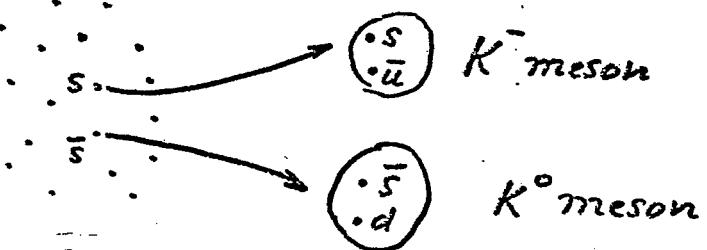
Signals for plasma formation (in pictures)

i) electromagnetic radiation emitted by plasma

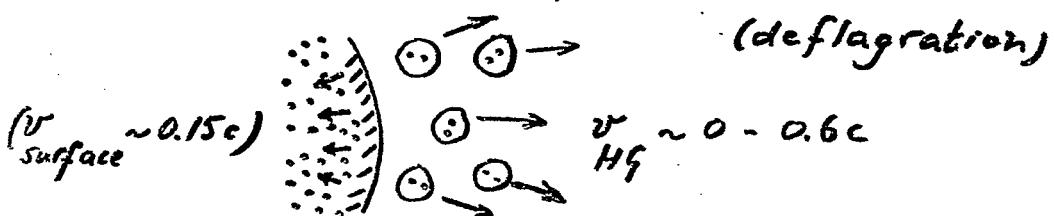


(escapes with negligible absorption)

ii) strange quark pairs formed in plasma, emitted as strange hadrons



iii) sudden liberation of latent heat at transition

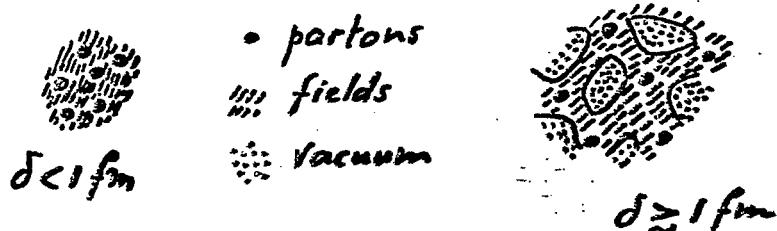


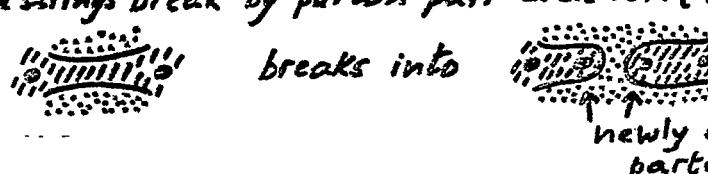
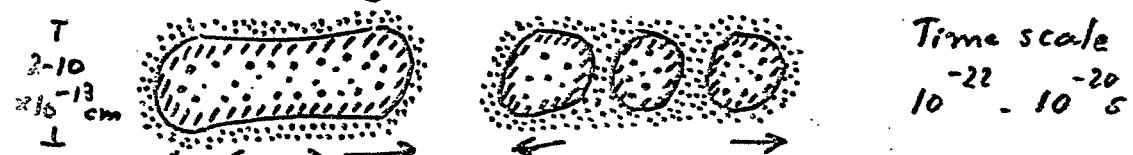
Plasma  $\longrightarrow$  Hadrons  
 high energy contents  $\rightarrow$  high transverse motion  
 " entropy "  $\rightarrow$  " multiplicity"

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Proposed scenario for Hadronization of expanding droplet of QG Plasma

- $\delta$ : nearest neighbour distance between QCD "partons" (= quarks, antiquarks, gluons)
- For  $\delta < 1 \text{ fm}$  colour confining fields (i.e. QCD gauge fields) cover all space in plasma.
- When  $\delta \gtrsim 1 \text{ fm}$  the fields collapse into strings (=flux tubes) separated by vacuum (non-perturb. QCD vacuum)

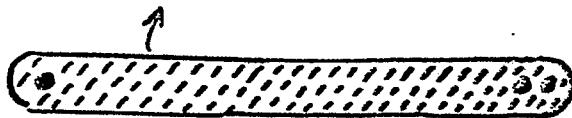


- Some strings break by parton pair creation (light quarks)
 
- A small expanding droplet breaks into smaller ones, and string tension stops expansion of the pieces
 

- Non-expanding plasma droplet can "deflagrate" by emitting a flow of hadron gas into vacuum.  
(In larger expanding droplets, bubbles of vacuum would form)

Heavy hadrons of high spin:  $J = \alpha' m^2 + c^t$

$$\alpha' \approx 1 \text{ GeV}^{-2}$$



- String of length  $L$ , diameter  $D \approx 1 \text{ fm} = \frac{1}{200 \text{ MeV}}$
- Mass of non-rotating string  $L\sigma$

$$\sigma = \frac{\pi}{4} D^2 T_{\text{oo}}$$

- Rotating string, endpoints velocity  $c=1$

$$\text{mass: } \frac{\pi}{2} L\sigma = m$$

$$\text{spin: } \frac{\pi}{8} L^2 \sigma = J$$

$$J = \frac{m^2}{2\pi\sigma}, \quad \alpha' = \frac{1}{2\pi\sigma}$$

Using  $D=1 \text{ fm}$ ,  $\alpha' = 1 \text{ GeV}^{-2}$  one finds

$$\sigma = 0.16 \text{ GeV}^2, \quad T_{\text{oo}}^S = 10^{-2} \text{ GeV}^4 = 1.25 \frac{\text{GeV}}{\text{fm}^3}$$

An ideal gas of gluons and light quarks ( $u, d$ ) at  $T_{\text{oo}} = 150 \text{ MeV}$  would have  $T_{\text{oo}}^f \approx 1 \text{ GeV/fm}^3$ .

# Colour fields in expanding plasma

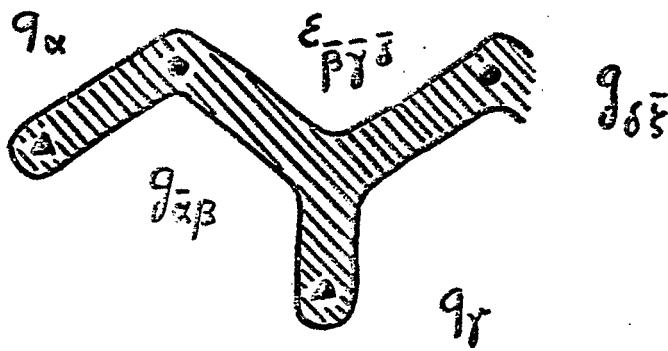
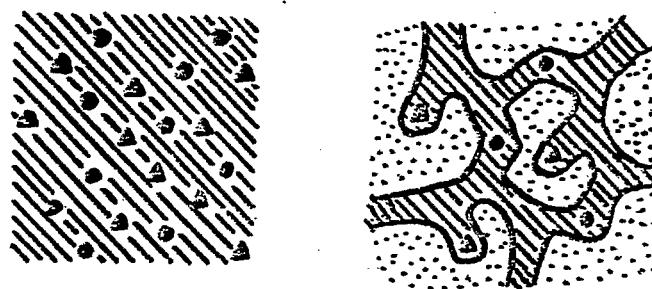
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- quarks and antiquarks
- gluons
- /// colour fields

QCD vacuum

$$1 \text{ fm} = 10^{-13} \text{ cm}$$

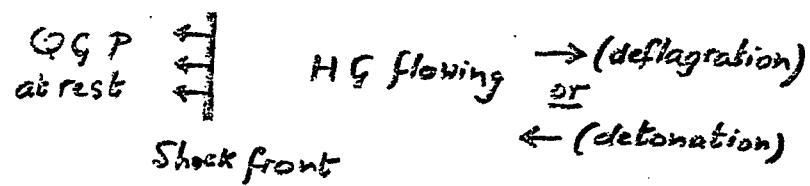
$$1 \text{ fm}$$



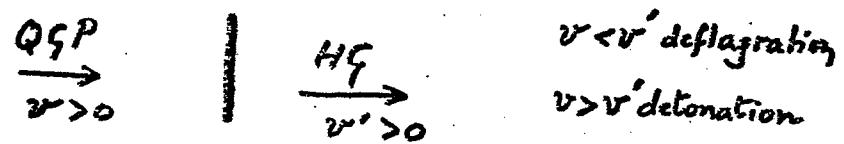
Can the Quark Gluon Plasma deflagrate?

[12]

Can it liberate its latent heat by emitting a flow of Hadron Gas?  
Picture in QGP rest frame:



Picture in rest frame of shock front:



Question: is this possible with increase of entropy?  
[If entropy would decrease, inverse process ( $v < 0$ ,  $v' < 0$ )  
could take place: compression of HG into QGP]

Energy conservation:  $\frac{v(\epsilon+p)}{1-v^2} = \frac{v'(\epsilon'+p')}{1-v'^2} \rightarrow vv' \geq 0$

$(QGP) \quad (HG)$   
 $\epsilon+p > \epsilon'+p' \rightarrow |v| < |v'|$  i.e. deflagration ( $0 < v < v'$ )  
or compression ( $0 < -v < -v'$ )

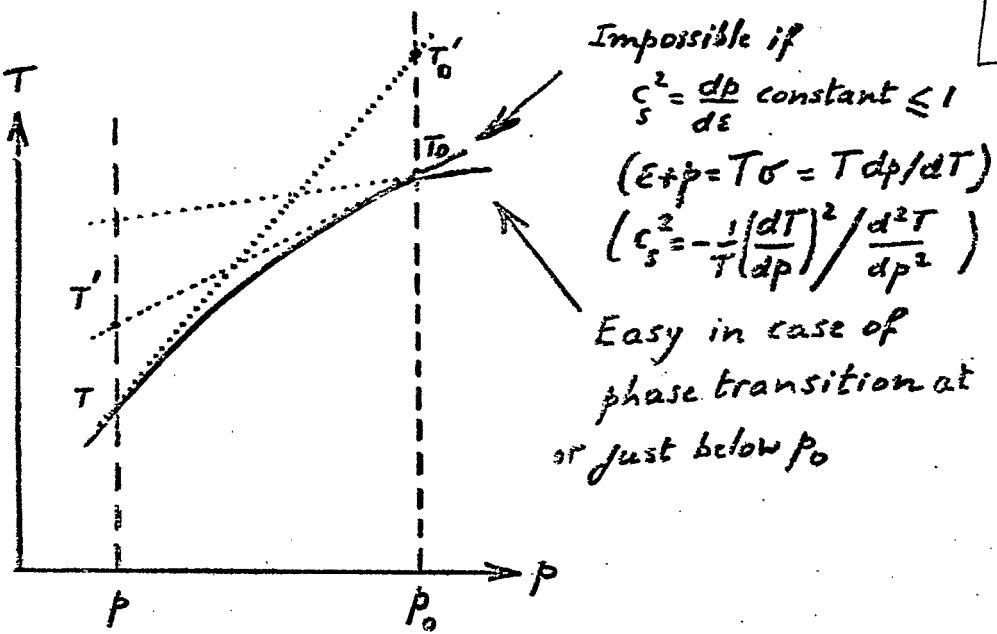
+ Momentum cons.:  $\frac{v^2}{1-v^2} = K \frac{\epsilon'+p'}{\epsilon+p}$ ,  $\frac{v'^2}{1-v'^2} = K \frac{\epsilon+p'}{\epsilon'+p}$ ,  $K = \frac{p-p'}{\epsilon-p-\epsilon'+p'}$

Baryon nb cons.:  $v^2 \frac{\epsilon+p}{\epsilon+p} = v'^2 \frac{\epsilon+p'}{\epsilon'+p'}$

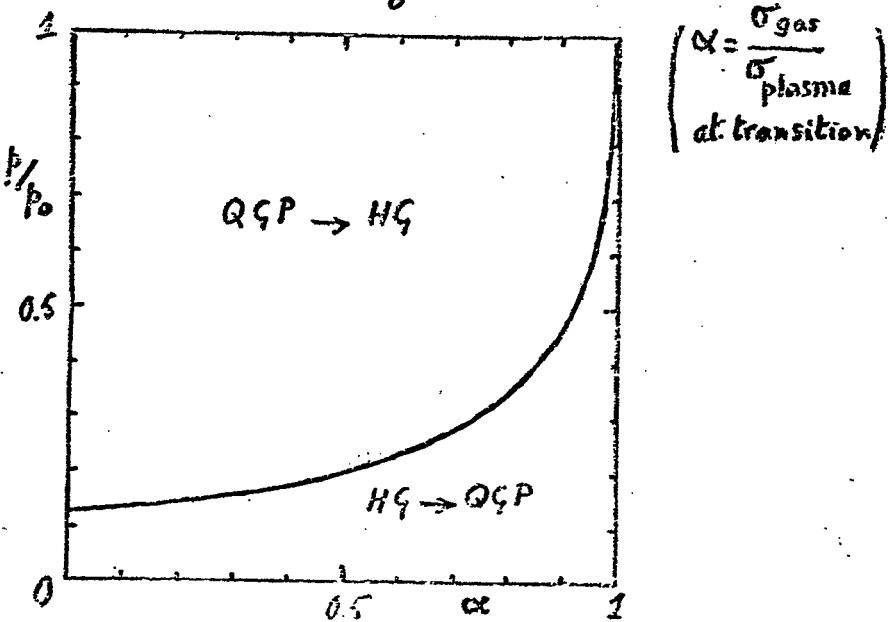
Entropy increase:  $\left\{ \begin{array}{l} v \text{ and } v' \neq 0: \frac{\sigma}{\sigma} \leq \frac{\sigma'}{\sigma}, \\ v=v'=0: \frac{\sigma^2}{\sigma'^2} \leq \frac{(\epsilon+p)(\epsilon+p')}{(\epsilon'+p)(\epsilon'+p')} \leftrightarrow \frac{T'^2}{T^2} \leq \frac{\epsilon+p'}{\epsilon+p} \frac{\epsilon'+p'}{\epsilon'+p} \end{array} \right.$

Graphical treatment of  $v=v'=0$  case: entropy increase  $T T'_0 \leq T' T_0$

[13]

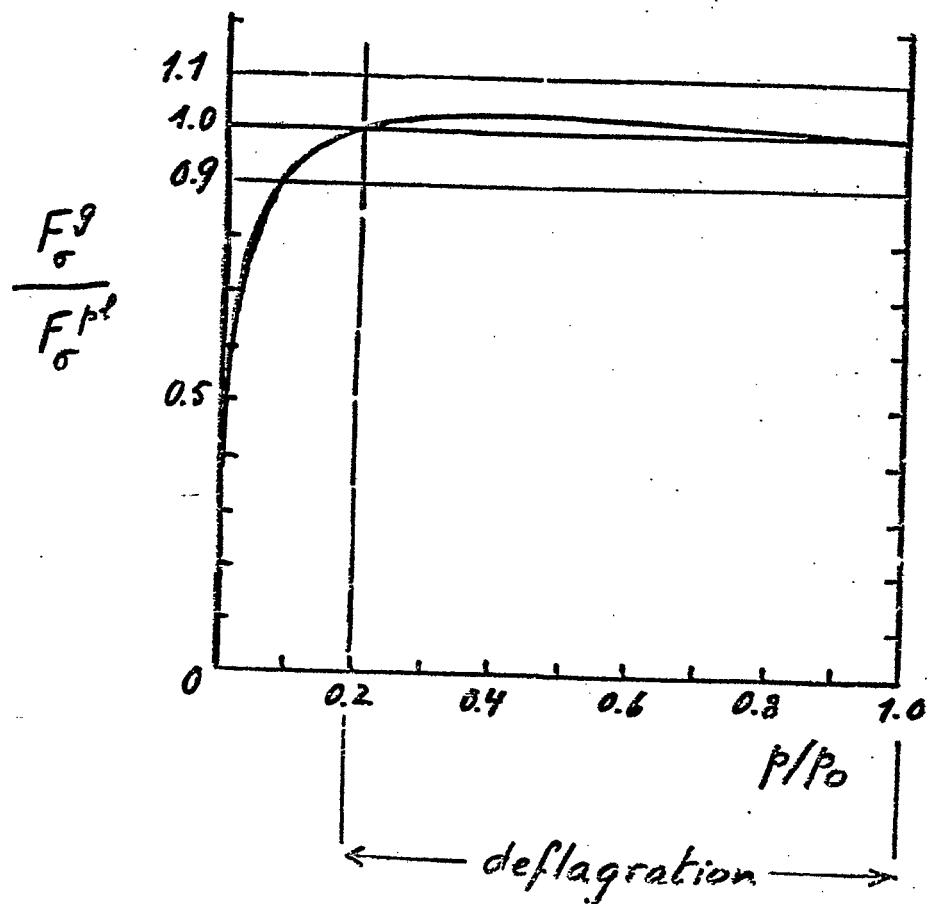


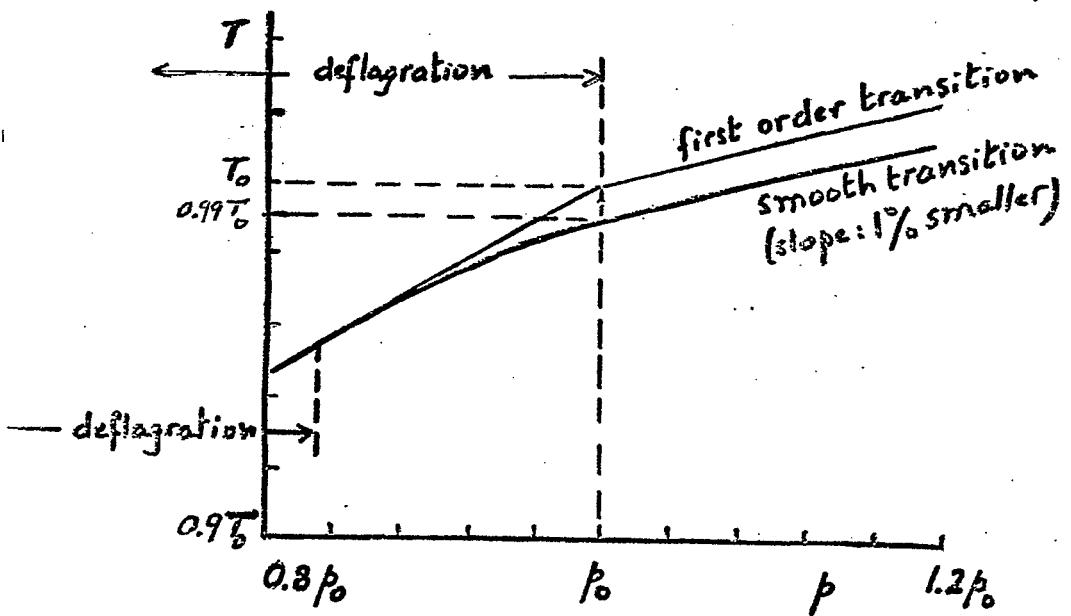
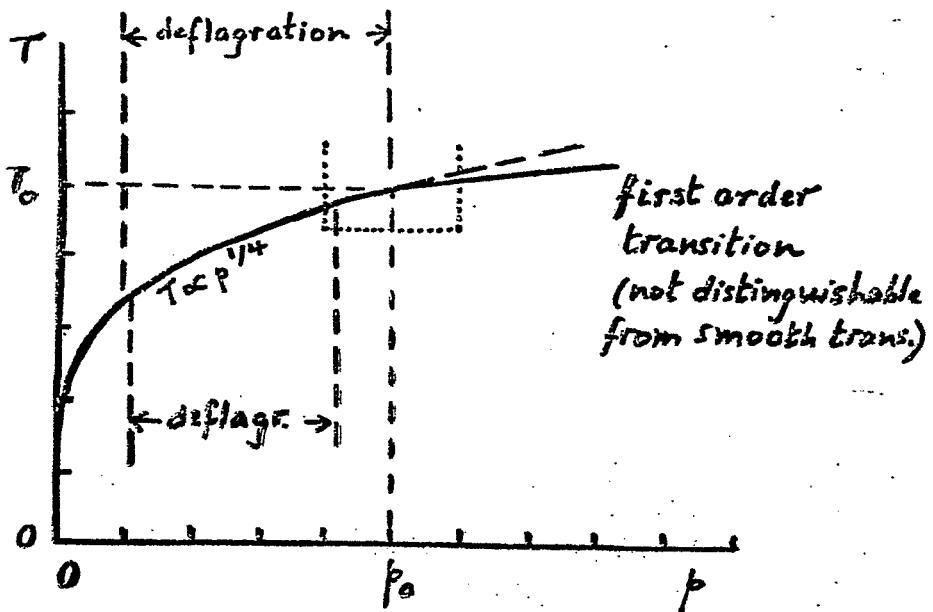
Example:  $T \propto p^{\frac{n}{n-1}}$ ,  $n=ct$  for  $p \leq p < p_0$ ,  $c_s^2 = \frac{n}{1-n}$ ,  $n \leq \frac{1}{2}$ . Take  $n = \frac{1}{4}$   
Phase transition at  $p$ :  $\left. \frac{dT}{dp} \right|_{p=0} = \alpha \left. \frac{dT}{dp} \right|_{p=p_0}$ ,  $\alpha < \infty$ . [ideal masses]

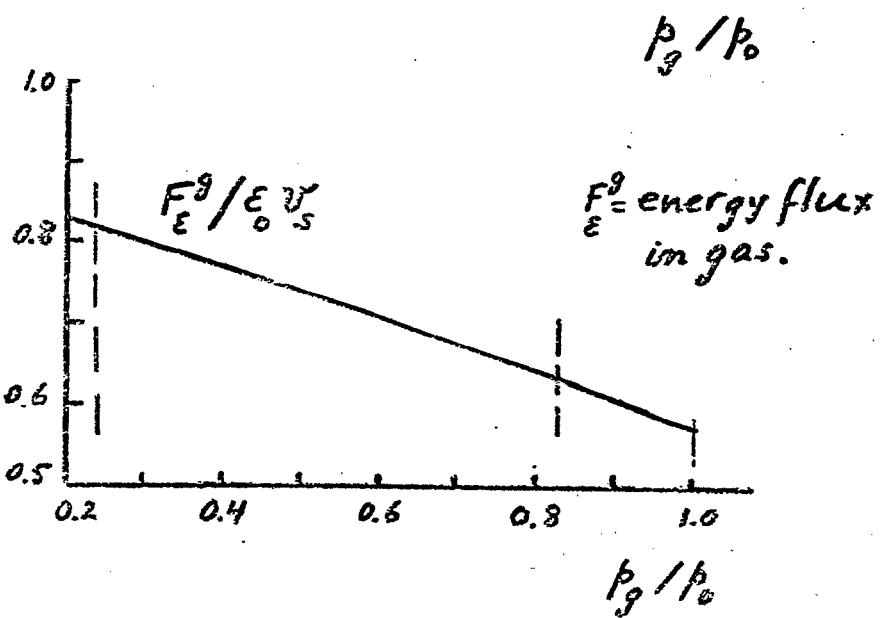
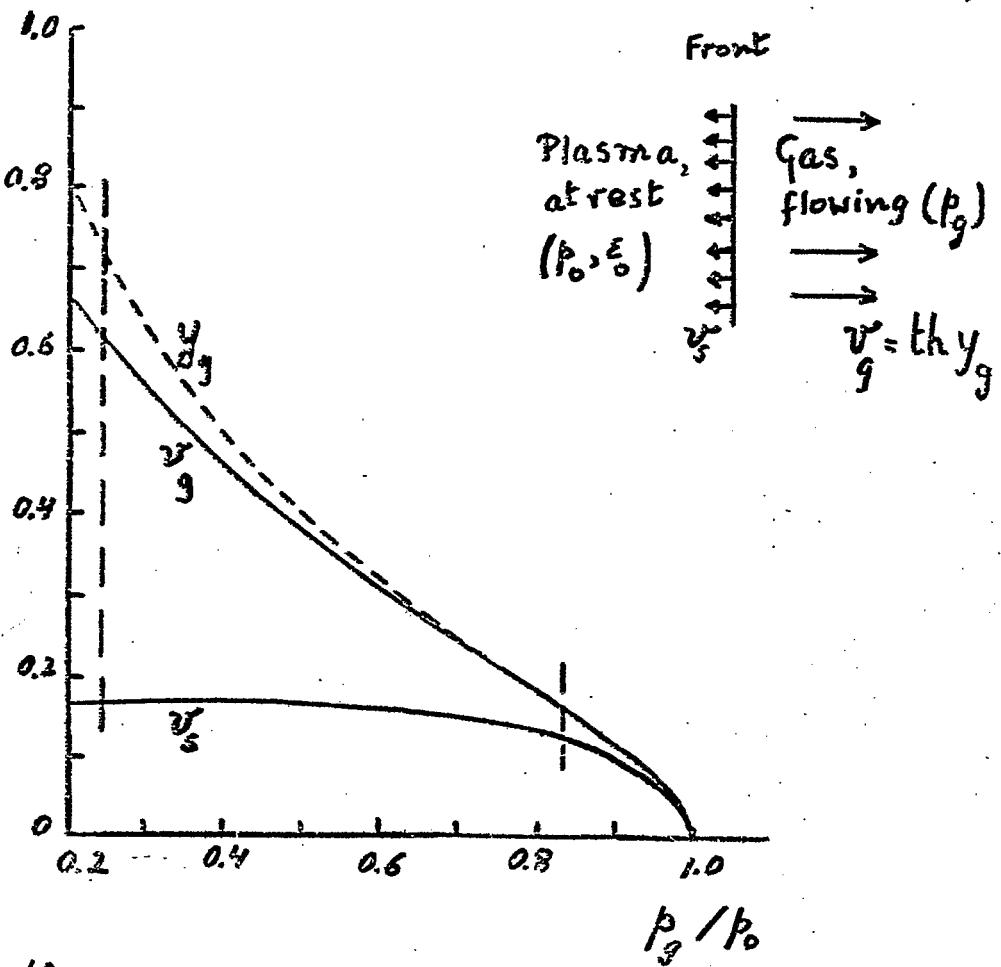


$F_g^g$ : entropy flux in gas  
 $F_p^{pl}$ : entropy flux in plasma } in rest frame of front.

First order trans.,  $\alpha = \sigma_g / \sigma_{pl} = 0.5$  at transition.







$$\text{Discussion of energy flux } F_{\epsilon}^g = \left(\frac{\epsilon + p_g}{g}\right) v_s \left(1 - v_s^2\right)^{-1} \quad [17]$$

- Order of magnitude  $\frac{3}{4} \times \frac{1}{6} \times \epsilon_0 = \frac{1}{8} \epsilon_0$

$$\frac{F_{\epsilon}^g}{\epsilon v_s}$$

- What would it be if all quarks, antiquarks and gluons flying toward surface would cross it without energy loss?

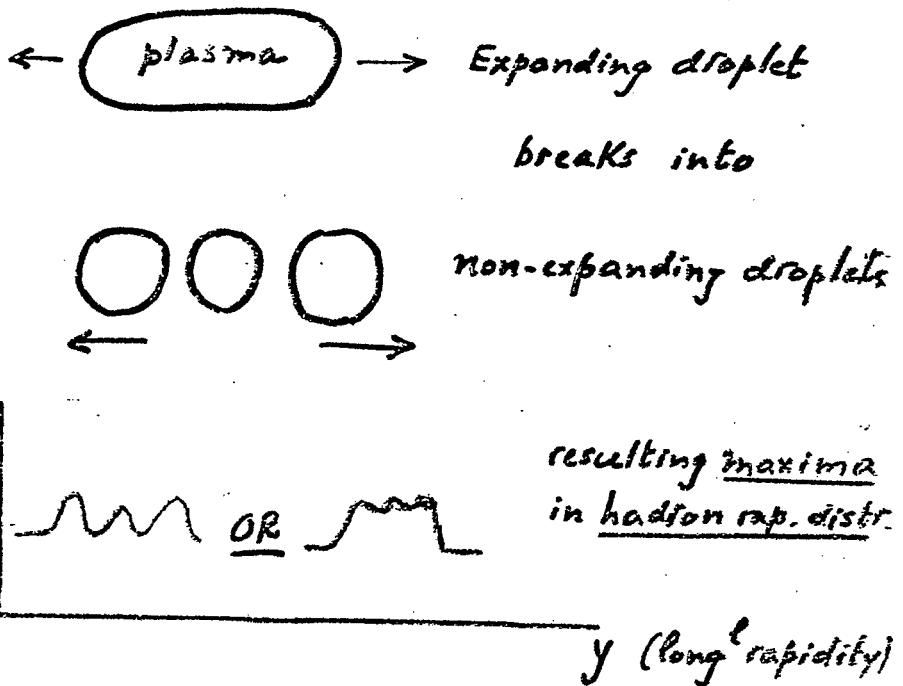
$$\frac{1}{2} \times \frac{1}{2} \times \epsilon_0 = \frac{1}{4} \epsilon_0$$

$$\frac{1}{2} \cos \theta$$

half of the quanta

- Experimental consequences.

If plasma droplet (possibly after breaking up) hadronizes by deflagration, resulting rapidity distribution of hadrons should show maxima at rapidities of droplets



Expected width of maxima  $\sim 1 \div 1.5$

- The other signals for plasma formation (direct dileptons and photons, strange hadrons) should be concentrated in these maxima.
- Hadrons from plasma should have  $p_T$  somewhat larger than normal, with broad but fluctuating azimuthal distribution.

THE HADRONIC PHASE TRANSITION  
IN THE EARLY UNIVERSE.

L. VAN HOVE

A Colloquium Lecture at  
Brookhaven National Laboratory  
(November 1984)

<u>"Age of Universe"</u>	<u>Temperature</u>	<u>Matter in the universe</u>
$\sim 2 \times 10^{10} \text{ y}$ $(\approx 6 \times 10^{17} \text{ s})$	$3^\circ \text{K}$ $(\approx 2.6 \times 10^{-4} \text{ eV})$	$0.4 \text{ photon/mm}^3, 0.1-1 \text{ nucleon/mm}^3, \frac{n}{p} \approx \frac{1}{8}, \frac{\rho}{p} = 1, \dots$ [ a few neutrinos / mm <sup>3</sup> (on average) ] + "dark" matter?
$\sim 6 \times 10^5 \text{ y}$	$\sim 3000^\circ \text{K}$	$\uparrow$ photons decoupled from other matter, atoms exist $\uparrow$ $\downarrow$ photons in equil. with plasma of nuclei and electrons $\downarrow$ $10^9 \text{ photons/mm}^3$
$\sim 3 \text{ min}$ $(= 180 \text{ s})$	$\sim 0.1 \text{ MeV}$	$\uparrow$ neutrons bound in nuclei $\uparrow$ $\downarrow$ protons and neutrons form gas $\downarrow$ Energy formation
$\sim 15 \text{ s}$	$\sim 0.3 \text{ MeV}$	$\uparrow$ no electron-positron pairs $\uparrow$ $\downarrow$ such pairs in equil. with photons and protons $\downarrow$
$\sim 1 \text{ s}$	$\sim 1 \text{ MeV}$	$\uparrow$ neutrinos decoupled from other matter $\uparrow$ $\downarrow$ neutrinos in equil. with other matter $\downarrow$
$\sim 10^{-3} \text{ s}$	$\sim 30 \text{ MeV}$	$\uparrow$ no hadrons except protons and neutrons $\uparrow$ $\downarrow$ unstable hadrons $\downarrow$ Limit of "conventional" physics.

"The first three minutes"  
by S. Weinberg, Basic Physics

"Known" matter is mostly leptonic  
but there is evidence for "dark" matter!

Picture based on "new" physics and increasingly speculative!!

$\sim 5 \times 10^{-5} s$	$\sim 150 \text{ MeV}$	$\uparrow$ hadrons form gas $\uparrow$	Hadronic phase transition (1) (large decrease of energy, and entropy density)
$\sim 2 \times 10^{-5} s$	$\sim 250 \text{ MeV}$	$\downarrow$ hadronic matter coalesces in <u>quark-gluon plasma</u>	$\gg 10^{36}$ particles/ $\text{mm}^3$ $= 1 \text{ part.}/\text{fm}^3 \propto T^3$ .
$\sim 10^{-8} s$	$\sim 10 \text{ GeV}$ ( $= 10^4 \text{ MeV}$ )	$\uparrow$ electromagn. $\uparrow$ weak $\uparrow$ strong } interactions, photons $\uparrow$	Electroweak phase transition (2)
$\sim 10^{-12} s$	$\sim 10^3 \text{ GeV}$	$\downarrow$ electro-weak interactions merge, photons mixed with massless $Z^0$ , $W^\pm$ massless	
$\sim 10^{-32} s$	$\sim 10^{13} \text{ GeV}$	$\uparrow$ electroweak } interactions, gluons $\uparrow$ $\uparrow$ strong	Electro nuclear phase transition (3) (origin of baryon asymmetry? monopoles, strings, ...)
$\sim 10^{-36} s$	$\sim 10^{15} \text{ GeV}$	$\downarrow$ electro-nuclear interactions merge, gluons mixed with other particles	
$\sim 10^{-42} s$	$\sim 10^{18} \text{ GeV}$		Other type of expansion? Creation of matter & entropy out of coherent field energy?
$\sim 0$ (by definition)	$\sim 10^{19} \text{ GeV}$ (Planck energy)	FRONTIER OF IGNORANCE (quantum gravity?)	
<u>Experimental tests:</u> ① ultrarelativistic ion collisions, ② colliders [ $p\bar{p}$ and $e^+e^-$ ], ③ { proton decay expts, heavy monopole searches			

"Known" matter/hadronic  
Nature of "dark" matter?

### III

Basic equations for expanding Universe (homogeneous isotropic model)

- Metric:  $ds^2 = dt^2 - d\ell^2$ ,  $d\ell^2 = a^2(t) \left[ \frac{dr^2}{1-Kr^2} + r^2(d\theta^2 + \sin^2\theta d\phi^2) \right]$

$r, \theta, \phi$ : space coordinates of co-moving observers

$t$ : proper time of co-moving observers

$d\ell$ : space distance,  $a(t)$  scale parameter

$K$ : curvature constant ( $K=0$  flat space, open

$K<0$  hyperbolic space, open } if  $S^3$   
comm.

$K>0$  spherical space, closed)

- Einstein equations for gravitation give:

$$\left( \frac{d}{dt} \frac{da}{dt} \right)^2 = \frac{8\pi}{3} G\rho - \frac{K}{a^2} + \Lambda$$

↑                           ↑ cosmological constant

gravit. attraction term,  $G$ =Newton constant,

$\rho$ =non-gravitational energy density

"kinetic energy term"       $G^{1/2} = 1.6 \times 10^{-33} \text{ cm.} (= \frac{t}{P}) = 0.54 \times 10^{-43} \text{ sec.} (= t_p)$

$$G^{-1/2} = 1.2 \times 10^{18} \text{ GeV} = 2.2 \times 10^{-5} \text{ g} (= M_p)$$

( $P$  for Planck,  $\hbar=c=1$ )

- Energy conservation for co-moving volume  $V$ ,

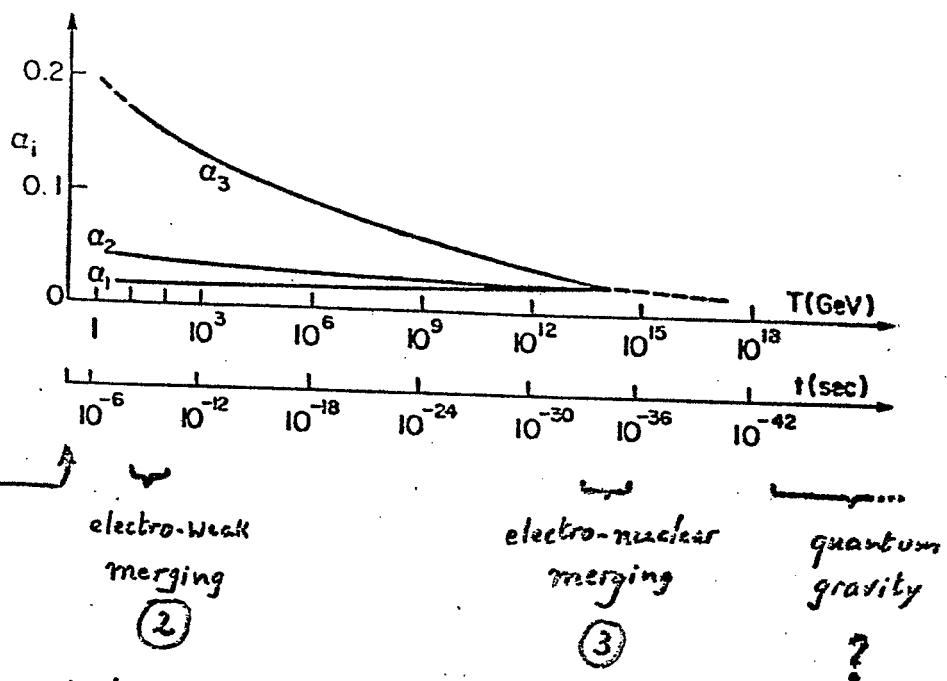
$$\frac{d}{dt}(\rho V) + p \frac{dV}{dt} = 0 \Rightarrow \frac{d}{dt}(\rho a^3) + p \frac{da^3}{dt} = 0$$

[ If  $\rho$  dominated by non-relativistic matter,  $p \ll \rho$ ,  $\rho a^3 = \text{constant}$  ]

[ " " " ultra- " " "  $p = \rho/3$ ,  $\rho a^4 = \text{constant}$  ]

(non-relativistic:  $5 \times 10^5 \text{ yr} \lesssim t \lesssim 1.5 - 2 \times 10^{10} \text{ yr} = \text{today}$ ,  
ultra- " before )

At every time, state of matter is characterized by the particle species present and their equation of state.



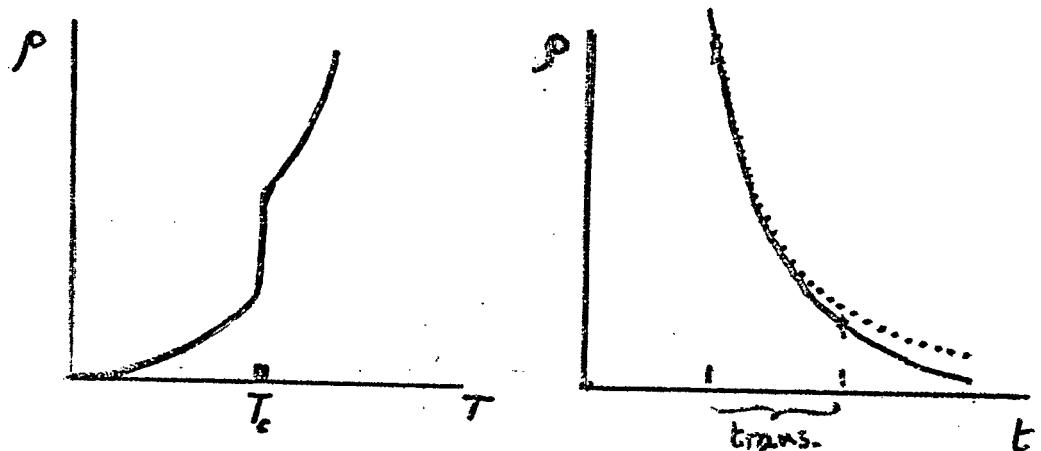
$\left[ \begin{array}{l} \text{electro-Weak} = \text{electromagnetic} + \text{weak interactions} \\ \text{electro-nuclear} = \text{electro-weak} + \text{strong interactions} \end{array} \right]$

$\alpha_i = \frac{g_i^2}{4\pi}$ ,  $g_i$ : coupling strength for  
energy-momentum transfers  $\sim T$

$\left[ \begin{array}{l} \alpha_1, \alpha_2 : \text{electro-Weak interactions} \\ - \alpha_3 : \text{strong interaction} \end{array} \right]$

Hadronic phase transition in the  
expansion of the Universe.

IV



$$\left(\frac{d}{dt} \frac{da}{dt}\right)^2 = \frac{8\pi G}{3} \rho \quad \frac{d}{dt} (\rho a^3) + \rho \frac{da^3}{dt} = 0$$

$\dot{\rho} \approx \rho_c$  constant during transition ]  
 $(\rho + \rho_c) a^3 \approx \dots$  [<sup>1st</sup> order trans.  
assumed for simplicity]

$$\frac{dp}{dt} = - (24\pi G)^{1/2} \rho^{1/2} (\rho + \rho_c)$$

$$\rho \approx \rho_c \left( \frac{\sqrt{\rho_1} - \sqrt{\rho_c} \tan \gamma(t-t_1)}{\sqrt{\rho_c} + \sqrt{\rho_1} \tan \gamma(t-t_1)} \right)^2$$

$\gamma = (6\pi G \rho_c)^{1/2}$ ,  $t_1$  = beginning of transition,  $t_1 \sim 10^{-5} s = 3 \text{ km}$

when  $\rho(t_1) = \rho_1 = \rho [QGP + \text{leptons} + \gamma's]$  at  $T_c, \rho_c$ .

Transition ends at  $t_2$  when  $\rho(t_2) = \rho_2 = \rho [H\zeta + \text{lept.} + \gamma]$  at  $T_c, \rho_c$ .

$$\rho_2 \sim 0.5 - 0.7 \rho_1, t_2 \sim 1.2 - 1.5 t_1.$$

At  $t = t_1 + \delta t$ ,  $H\zeta$  bubbles separated by  $t_1^{1/3} (1 \text{ fm})^{2/3} \sim 10^{-7} \text{ cm}$ .  
After percolation space structure may reach  $t_2^{2/3} (1 \text{ fm})^{1/3} \sim 1 \text{ cm}$

Most probably, no trace is left today of the hadronic phase transition in the early Universe

But one cannot exclude some exotic possibilities:

1) Black hole formation (Schramm et al.)

If QGP regions of size  $\sim t$ , surrounded by larger HG regions would form, some of them could collapse into black holes of mass  $\lesssim 10^{58} \text{ GeV}$  and radius  $\lesssim 10 \text{ km}$  ( $M_\odot = 2 \times 10^{33} g \approx 10^{57} \text{ GeV}$ )

The black holes could later contribute to galaxy formation.  
They would now form invisible matter

(H. Crawford & D.N. Schramm, Nature 298 (1982) 538;  
K. Freese, R. Price & D.N. Schramm, Ap.J. 275 (1983) 405)

For this to work, the  $\rho_{BH}$  in the black holes  
should now be  $\approx (10-100) \rho_{\text{visible}}$

2) Formation of "strange quark matter" (Witten):

a possible new form of stable matter containing u, d, s quarks, could also form invisible matter in the present Universe

(E. Witten, Phys. Rev. D 30 (1984) 272)

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There is something fascinating about Science -

One gets wholesale returns of conjecture

out of such trifling investment of fact!

(Mark Twain)

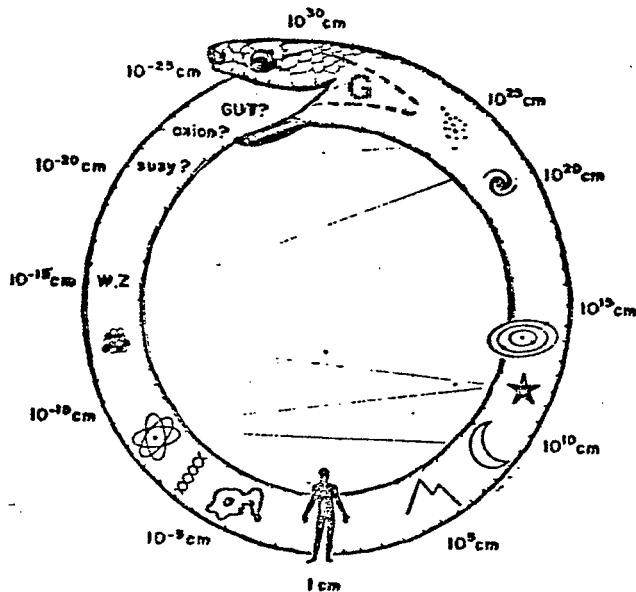


Fig. 7. New ideas in particle physics suggest further linkages between "micro" and "cosmic" scales (i.e. between left and right on this picture) in addition to the well-known ones between atomic and nuclear physics and the terrestrial and stellar scales. The ultimate unification will involve quantum gravity (the Planck scale) and cosmology. (from Primack and Blumenthal, Proc. 3rd Moriond Astrophysics Workshop (Reidel) p. 180 (1983)).